

A Centered Hydrogen Bond in Potassium Hydrogen Chloromaleate: A Neutron Diffraction Structure Determination*

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The crystal structure of potassium hydrogen chloromaleate has been directly determined from three-dimensional neutron diffraction data by means of the Patterson synthesis. The nearly plane anion has a ring structure with a 2.40 Å O···H···O bridge in which the hydrogen atom is effectively centered.

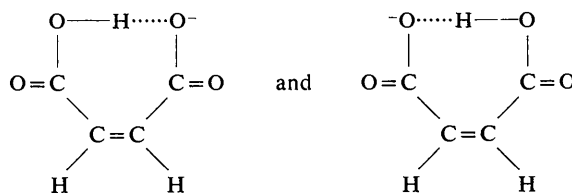
Introduction

Hydrogen bonds between oxygen atoms are known to occur with O···O distances ranging upward from about 2.4 Å. In the great majority of cases the hydrogen atom is situated unsymmetrically in the bridge at about 1 Å from one of the two oxygen atoms. The actual O—H distances and the fundamental O—H stretching frequencies show a pronounced correlation with the length of the hydrogen bond, such that the longest O—H distances and the lowest frequencies tend to occur in the shortest O···O bridges. The available data have been summarized by Bacon (1963), Pimentel & McClellan (1960), Rundle & Parasol (1952), and Lord & Merrifield (1953).

The trend of the variation in O—H distance with O···O distance suggests that the hydrogen atom may take a symmetric location in some of the shortest of the hydrogen bonds, those with O···O distances less than about 2.5 Å. A similar indication is derived from the behavior of the O—H stretching frequency, which becomes rapidly lower as the O···O distance approaches the neighborhood of 2.5 Å.

The most firmly established example of a centered hydrogen bond is the bifluoride ion F—H—F⁻, for which definitive evidence from thermodynamics (Westrum & Pitzer, 1949), spectroscopy (Mathieu & Couture-Mathieu, 1950), nuclear magnetic resonance (Waugh, Humphrey & Yost, 1953), and neutron diffraction (Peterson & Levy, 1952; McGaw & Ibers, 1963; Ibers, 1964) is available. Among O···H···O bonds, one, that in chromous acid, appears to be firmly established as centered (Hamilton & Ibers, 1963), and centering has been suggested in several other cases, among them potassium hydrogen maleate (Cardwell, Dunitz & Orgel, 1953). This crystal contains one of the shortest O···H···O bonds known. Its structure has been studied by both neutron (Peterson & Levy, 1958) and X-ray (Darlow & Cochran, 1961) diffraction. Both studies fixed the apparent position of the hydrogen atom at the center of a 2.44 Å intraionic O···H···O

bridge. The apparent centering of the hydrogen atom as given by the diffraction studies can be given two interpretations; namely, the hydrogen atom is truly localized at the center of the bond, or it is randomly distributed in one or the other of two equivalent sites displaced perhaps as much as 0.1 Å from the center, as indicated by the structural formulas:



The equivalence of the two carboxyl groups, as required by the presence of a mirror plane in the crystal, makes the two alternatives indistinguishable. If this equivalence is removed, as by substitution for one of the ethylenic hydrogen atoms, the two non-centered configurations become non-equivalent; and there is a chance of directly establishing the existence of the non-centered situation, should it occur. Accordingly, the crystal structure of the potassium salt of the chlorine-substituted acid has been studied. To our knowledge, this is the first study of a very short O···H···O bond in which the oxygen atoms are crystallographically and chemically non-equivalent.

Experimental

The space group and cell parameters of potassium hydrogen chloromaleate, already reported by Ellison, Levy & Peterson (1961), are as follows: *Pbcn*, $a = 15.815$ (standard deviation 0.015), $b = 10.928$ (0.006), $c = 7.707$ (0.005) Å, eight molecules in the unit cell. A crystal specimen with volume about 5 mm³, grown from aqueous solution, was used for this study.

A three-dimensional set of 1959 symmetrically non-equivalent intensities, complete to $\sin \theta/\lambda = 0.71 \text{ \AA}^{-1}$, was measured with the Oak Ridge automatic neutron diffractometer (Busing, Smith, Peterson & Levy, 1964) using a neutron wave length of 1.078 Å. Net counts were obtained in a manner already described (Brown

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& Levy, 1964). Absorption corrections were computed with a FORTRAN program written by Wehe, Busing & Levy (1962). The range of transmission factors was 0.844 to 0.899. Squared structure factor amplitudes were derived in the usual manner after averaging replicate and crystallographically equivalent intensities. Standard errors for use in least-squares weighting were estimated according to the expression (Brown & Levy, 1964)

$$\sigma^2(F^2) = K^2[G + (t_P/t_B)^2B + (0.03N)^2]$$

in which G is the gross peak count (including background), B the background count, N the net count, (t_P/t_B) the ratio of counting times of peak to background, and K the appropriate scaling constant. The third term within the brackets is an empirical one justified by previous experience with the instrument employed (Brown & Levy, 1964).

Solution of the structure

In preliminary trials it became apparent that the high symmetry of the space group, which results in a Patterson function containing the superposed interactions of eight molecules, makes direct interpretation excessively difficult, particularly in the absence of a pronounced 'heavy' atom (The largest neutron scattering factor, that for chlorine, is 1.7 times the mean. The total number of atoms in the unit cell is 96; of these, 72 are of approximately equal scattering power).

The procedure that proved successful was derived independently in this work. Although closely similar methods have been described by others (Hamilton, 1963; Mighell & Jacobson, 1963), a description in mathematical terms seems worthwhile here. The procedure consists of a systematic examination of the Patterson function through the use of the Buerger minimum function. This was accomplished by computation of the function

$$Q(\mathbf{r}) = \text{Min}_i[w_i P(\mathbf{r} - \mathbf{S}_i \mathbf{r})] \quad (1)$$

for values of the position vector \mathbf{r} ranging over the asymmetric unit of $Pbcn$. In the foregoing equation, \mathbf{S}_i is a symmetry operator of the space group, w_i is a weighting factor equal to the reciprocal of the number of operations \mathbf{S}_i which generate a given vector $\mathbf{r} - \mathbf{S}_i \mathbf{r}$, P is the Patterson function, and the symbol Min_i designates the minimum value of the quantity in brackets over the set of symmetry operations of the space group.

Of the several Patterson values compared in the minimum function of Q , one is $P(0, 0, 0)$, corresponding to the identity symmetry element; one is evaluated at a general point in Patterson space, corresponding to the inversion center (for centrosymmetric space groups); and the remainder are evaluated at points in the several Harker sections and lines. The latter components of $Q(\mathbf{r})$, considered separately, contain the same vector sets as the corresponding Buerger implication diagrams (Buerger, 1959), to which they are

closely related. They are, however, expressed in terms of continuous, rather than discrete variables. It seems appropriate to name $Q(\mathbf{r})$ the continuous multiple implication function.

The quantity $Q(\mathbf{r})$ will tend to assume positive maxima at those positions \mathbf{r} for which the corresponding Patterson-Harker values are all large and which are thus suitable candidates for atomic positions. These positions \mathbf{r} will occur in the function with respect to each of the eight centers of symmetry which are possible origins of $Pbcn$; this additional symmetry (over that of the unit cell) is reflected in the fact that for this space group $Q(\mathbf{r})$ is periodic in each dimension with a period of one-half that of the crystal lattice. This 'false symmetry' arises from the same considerations as do the ambiguities in implication diagram theory (Buerger, 1959).

In addition, positive regions may appear because of accidental coincidence of interaction vectors. For a neutron Patterson function, accidental cancellation of Patterson interaction vectors may also occur, since scattering amplitudes of both signs may be present.

In applying this method, we used a gradient-sharpened* Patterson function (Jacobson, Wunderlich & Lipscomb, 1961) with coefficients $F_{\frac{2}{3}}^2 = (\frac{1}{3} + \lambda^{-2} \sin^2 \theta) F_{\text{obs}}^2$. In the corresponding function $Q(\mathbf{r})$ there appeared one well-resolved sizeable maximum, the third largest on the diagram, with height intermediate between those expected for Cl-Cl and C-C. The second largest peak although less well resolved, was successfully used in subsequent analysis. The largest peak was obviously perturbed by overlap.

The resolved peak was made the basis for a multiple superposition diagram, in which the Buerger minimum function and the symmetry of the unit cell were again used. This function can be expressed in a fashion similar to that used for $Q(\mathbf{r})$:

$$S(\mathbf{r}) = \text{Min}_i[P(\mathbf{r} - \mathbf{S}_i \mathbf{r}_o)] \quad (2)$$

where \mathbf{r}_o is the tentatively assigned atomic position found from $Q(\mathbf{r})$. In $S(\mathbf{r})$, the eightfold ambiguity of origin is, of course, no longer present.

Comparison of the first S map with the Q map, taking suitable account of symmetry, suggested two other atomic positions, one of which corresponded to the second largest peak, and the other to a shoulder on the fourth largest peak of $Q(\mathbf{r})$. Calculations of $S(\mathbf{r})$ based on these vectors confirmed their self-consistency as atomic positions, and detailed comparison of the three superposition diagrams disclosed common features suggesting the expected shape of the chloromaleate ion. At this stage, the minimum of all three functions

* Because of a programming error, the Patterson function was at first sharpened only in the x direction. Nevertheless, the structure was successfully solved from this Patterson function and the subsequent Q and S maps. Fig. 1 is prepared from the correctly sharpened Patterson function. In retrospect, it appears that the structure could probably have been seen in the first S map, and certainly in the second, had the error not been made.

$S(r)$ was prepared; the resulting diagram displayed maxima assignable to all atoms except one potassium and the two hydrogen atoms. These peaks are absent because of the relatively small scattering amplitudes of K for neutrons, and because of the negative sign of interaction of H with the peaks chosen for superposition. This 24-fold minimum function contains no extra peaks large enough to be troublesome. A composite contour diagram is presented in Fig. 1.

Refinement of the structure

A preliminary Fourier synthesis of the scattering density with phases based on the atomic positions deduced from the Patterson function analysis revealed peaks corresponding to all atoms. The carbon and oxygen atoms are in general positions of the space group; the potassium atoms appear in special positions on diad axes parallel to b and are of two crystallogra-

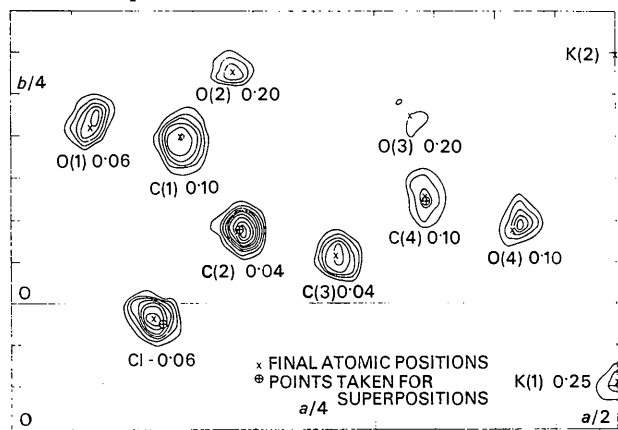


Fig. 1. Composite contour map of the multiple superposition diagram. The decimal fractions indicate the z coordinates of the levels from which the contoured portions were taken.

phically distinct kinds. The two hydrogen atoms in the asymmetric unit were indicated to be in general positions by the only two sizeable regions of negative density. A second Fourier synthesis followed, with phases based on atomic parameters derived from the first.

Following the Fourier syntheses, a sequence of least-squares refinements was carried out in which were varied successively (1) the overall scale factor and individual isotropic atomic thermal parameters, (2) the foregoing plus atomic position coordinates, (3) the scale factor, anisotropic temperature coefficients, and atomic position coordinates, and (4) the scattering amplitudes of chlorine, hydrogen, oxygen, and potassium along with those parameters listed under (3). In refinement (4), the four strongest reflections, which

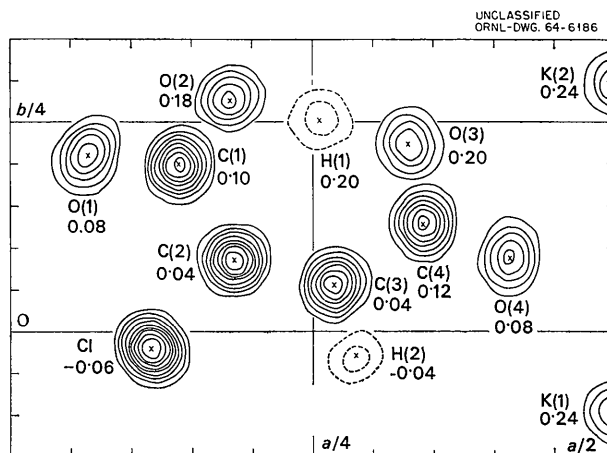


Fig. 2. Composite contour map of the final three-dimensional Fourier synthesis of the structure. The contour interval is $0.48 \text{ cm} \times 10^{-12} \text{ \AA}^{-3}$; zero contours are omitted, and negative contours are dotted. The decimal fractions indicate the z coordinates of the levels from which the contoured portions were taken. The crosses indicate final x and y coordinates.

Table 1. Parameters of the structure of potassium hydrogen chloromaleate

Least-squares standard errors are given in parentheses

Atom	Fractional position parameter $\times 10^4$			Thermal parameter* $\times 10^5$					
	x	y	z	B_{11}	B_{22}	B_{33}	B_{12}	B_{13}	B_{23}
Cl	1161(1)	-197(1)	-534(2)	265(4)	642(10)	1407(21)	-74(5)	-7(8)	-268(13)
C(1)	1394(1)	1988(1)	1041(2)	200(4)	524(11)	871(21)	64(6)	-15(9)	-59(14)
C(2)	1850(1)	860(1)	381(2)	181(4)	459(10)	868(22)	21(6)	25(8)	-22(14)
C(3)	2677(1)	575(1)	387(2)	189(4)	530(12)	1315(26)	57(6)	29(9)	-3(16)
C(4)	3414(1)	1281(2)	1108(2)	167(4)	701(13)	1323(29)	26(6)	-20(10)	188(17)
O(1)	638(1)	2100(2)	727(3)	226(6)	902(19)	1412(40)	173(9)	-82(12)	-302(23)
O(2)	1817(1)	2767(2)	1918(3)	293(7)	571(16)	1832(43)	54(9)	-76(15)	-360(21)
O(3)	3290(1)	2234(2)	2039(4)	245(7)	975(24)	2139(55)	-41(11)	-131(16)	-392(28)
O(4)	4131(1)	880(2)	787(3)	173(6)	963(21)	1984(49)	81(9)	21(13)	297(27)
H(1)	2558(3)	2522(4)	2032(6)	384(14)	868(31)	1932(68)	-66(16)	-97(29)	-332(41)
H(2)	2863(2)	-289(3)	-203(6)	346(13)	785(33)	2755(107)	132(16)	131(31)	-368(47)
K(1)	5000	-946(4)	2500	258(13)	473(31)	1099(69)	0	-73(24)	0
K(2)	5000	2973(4)	2500	191(12)	749(39)	1162(73)	0	80(23)	0
Scattering amplitudes, $\text{cm} \times 10^{-15}$				C	Cl	O	H	K	
				661	937(5)	586(3)	-367(3)	363(4)	

(Amplitudes of atoms of the same chemical element were constrained to be equal. That of carbon was not varied).

* Coefficients in the expression $B_{11}h^2 + B_{22}k^2 + B_{33}l^2 + 2B_{12}hk + 2B_{13}hl + 2B_{23}kl$.

were clearly subject to extinction error, were omitted. The value of $\sum w(F_o^2 - F_c^2)^2$ was minimized.

The final value of the standard deviation of fit

$$\sigma_1 = [\sum w(F_o^2 - F_c^2)^2 / (N - P)]^{1/2}$$

(N is the number of observations and P the number of adjusted parameters) is 1.015 (1.19 before adjusting scattering factors). The conventional discrepancy indices $R_x = \sum ||F_o|^x - |F_c|^x| / \sum |F_o|^x$ are summarized in the following table.

	A	B	C
$x=1$	0.122	0.123	0.079
$x=2$	0.095	0.092	0.081
Number of data:	1959	1955	1447

A—all data are included.

B—data which are affected by extinction are omitted.

C—data which are affected by extinction or have $F^2 < \sigma$ are omitted.

Although the values of the discrepancy factors are not especially low by modern standards, the close approach of the standard deviation of fit to unity indicates that the parameters have been refined as fully as the data warrant. This conclusion is confirmed by an $F_o - F_c$ Fourier synthesis, which shows no interpretable

features, the extreme density values being $+0.166$ and $-0.162 \text{ cm}^3 \times 10^{-12}$ per \AA^3 . For comparison, a hydrogen peak has density -1.30 and a potassium peak $+1.95$.

The precision of the parameters is also quite satisfactory: the mean positional parameter standard errors range from 0.0013 \AA for Cl to 0.0043 \AA for H and K.

The final parameters and the least-squares standard errors are listed in Table 1. The observed and calculated values of F^2 are shown in Table 2. A composite Fourier synthesis of the structure is illustrated in Fig. 2.

Results

Structure of the ion

As expected, the hydrogen chloromaleate ion has a closed-ring structure, analogous to that of the unsubstituted hydrogen maleate ion. This structure is illustrated in Figs. 3 and 4 by means of stereoscopic drawings prepared on a computer-controlled plotter, which was programmed by our colleague C. K. Johnson (Johnson, 1965). Bond distances and angles are shown in Fig. 5.

The internal hydrogen bond distance $O(2) \cdots O(3)$ has the unusually small value 2.403 \AA , slightly shorter

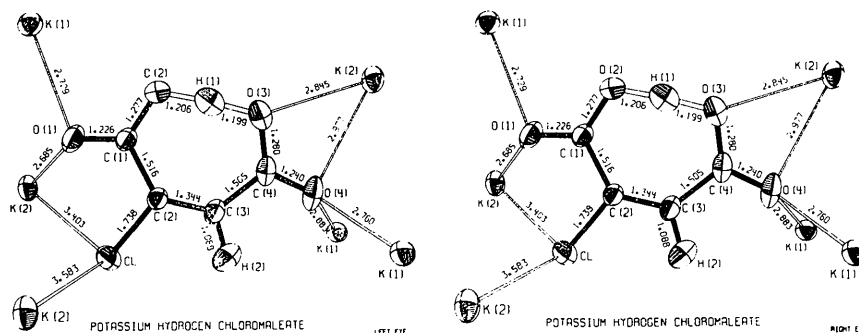


Fig. 3. A stereoscopic pair of drawings of the chloromaleate ion viewed in its own plane. The atoms are represented by ellipsoids of thermal displacement including 50% probability. Figs. 3 and 4 are best viewed by means of an inexpensive stereoscope. (The addresses of manufacturers of suitable stereoscopes may be obtained from the authors on request).

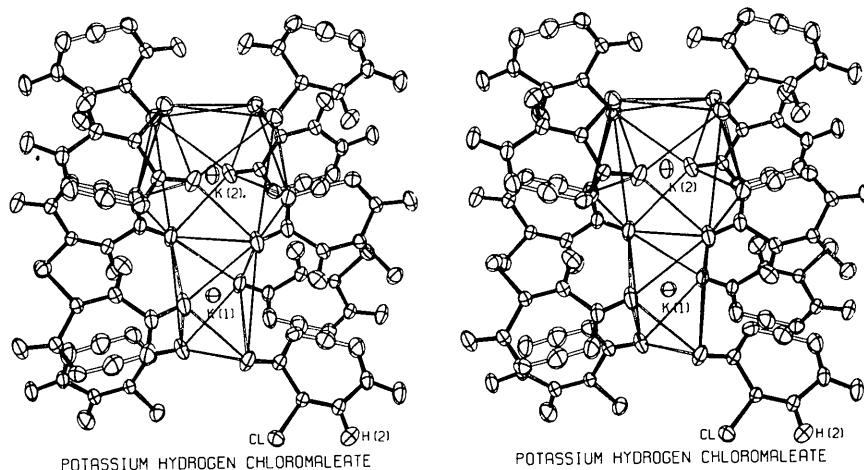


Fig. 4. Stereoscopic drawings of the coordination about potassium ions, viewed nearly along the c axis.

Table 2 (cont.)

Table with multiple columns of numerical data, organized into groups by row headers (e.g., H 1, H 2, H 3, H 4, H 5, H 6, H 7, H 8, H 9, H 10, H 11, H 12, H 13, H 14, H 15, H 16, H 17, H 18, H 19, H 20, H 21, H 22, H 23, H 24, H 25, H 26, H 27, H 28, H 29, H 30, H 31, H 32, H 33, H 34, H 35, H 36, H 37, H 38, H 39, H 40, H 41, H 42, H 43, H 44, H 45, H 46, H 47, H 48, H 49, H 50, H 51, H 52, H 53, H 54, H 55, H 56, H 57, H 58, H 59, H 60, H 61, H 62, H 63, H 64, H 65, H 66, H 67, H 68, H 69, H 70, H 71, H 72, H 73, H 74, H 75, H 76, H 77, H 78, H 79, H 80, H 81, H 82, H 83, H 84, H 85, H 86, H 87, H 88, H 89, H 90, H 91, H 92, H 93, H 94, H 95, H 96, H 97, H 98, H 99, H 100). Each group contains several rows of numbers, some with sub-headers like 'CALC OBS SIG'.

than the value 2.437 Å reported for the unsubstituted compound. The hydrogen atom appears to be, within experimental error, equidistant from the two oxygen atoms. Assessment of the full significance of this apparent centering of the hydrogen bond involves consideration of the thermal motion of the crystal, and discussion of this matter will be deferred to a later section.

The chloromaleate ion as a whole is not quite plane: calculation of the least-squares best plane by the method of Hamilton (1961) shows the average deviation of the 11 atoms from the plane to be 0.076 Å, with extremes to the two sides of 0.19 Å for O(3) and 0.16 Å for O(1). These deviations are shown in Fig. 5. The ethylenic group, however, is quite closely plane; the extreme deviations from a least-squares best plane through the six atoms of this group are 0.03 Å for H(2) and 0.01 Å for C(4). Similarly, there are but small departures from planarity at each carbon atom: the largest of the six dihedral angles formed by the 4 planes that are defined by C(4) and its three nearest neighbors is 1.6°; the corresponding quantities for the others are 0.9° for C(1), 0.6° for C(2) and 0.1° for C(3) (standard errors are 0.3°).

The largest deviations from planarity of the ion correspond to rotations of the carboxyl groups about the C–C single bonds. As measures of these rotations, we give the conformation angles* C(3)–C(2)–C(1)–O(2), $-8.5^\circ \pm 0.3^\circ$, and C(2)–C(3)–C(4)–O(3), $+9.0^\circ \pm 0.3^\circ$. For comparison, the conformation angle about the double bond, C(1)–C(2)–C(3)–C(4), is $1.9^\circ \pm 0.3^\circ$. It is somewhat puzzling that the carboxyl groups are rotated in the sense that displaces O(2) and O(3) to the same side of the mean plane of the ion.

All of the bond lengths are close to expected values. The dimensions of the two carboxyl groups are almost identical, indicating that the bonding is similar.

The bond lengths in the hydrogen chloromaleate ion are very similar to those reported (Darlow & Cochran, 1961) for the unsubstituted ion. The greatest difference, that in the length of the C(3)–H(2) bond, 1.088 ± 0.004 Å (this study) *vs* 0.96 ± 0.02 Å, may be attributed to a systematic discrepancy which frequently occurs in the location of hydrogen atoms by X-ray diffraction. Other differences are in the length of the C(1)–C(2) single bond adjacent to the chlorine atom, which is greater in the present case by 0.018 ± 0.004 Å, and in the C(1)–O(2) ··· O(3) angle, greater by $2.2 \pm 0.3^\circ$. The foregoing errors are the combined precision measures from both determinations. As in the unsubstituted ion, both of the C–C–C angles show considerable strain (Darlow, 1961), and the considerations of strain energy presented there are also applicable to the hydrogen chloromaleate ion.

* Conformation angle $A-B-C-D$ is defined to be the azimuth of CD with respect to BA about the vector BC as polar axis of a right-handed system; that is, the clockwise angle from the projection of BA to the projection of CD viewed in the direction BC . An equivalent definition was given by Brown & Levy (1963).

Ionic packing

Fig. 3 shows the ionic contacts between one hydrogen chloromaleate ion and its potassium ion neighbors. It is seen that three of the four oxygen atoms and the chlorine atom have potassium ion neighbors at significant contact distances. The absence of an ionic contact with O(2) provides an environmental dissimilarity between the two carboxyl groups which makes their dimensional similarity all the more noteworthy.

The overall packing arrangement can best be appreciated by examining the environment of the potassium ions, which occur in infinite sheets parallel to the $b-c$ planes of the crystal, separated by the distance $a/2$. One type, K(1), is surrounded by six oxygen atoms at the vertices of a distorted octahedron possessing a diad axis of symmetry. The other, K(2), is surrounded by six oxygen atoms and four chlorine atoms at the vertices of a 14-sided polyhedron also possessing a diad axis. One polyhedron of each type is illustrated in Fig. 4, which shows the structure viewed nearly along

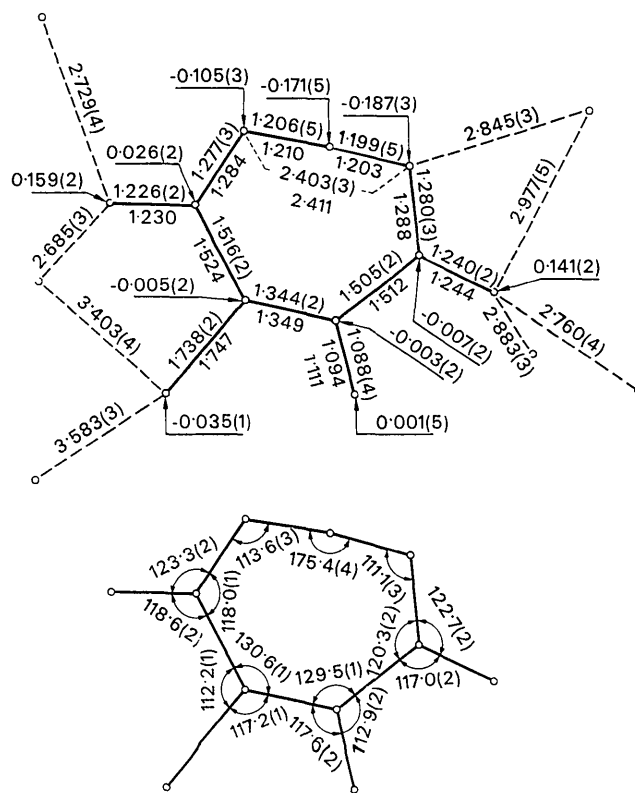


Fig. 5. Distances and angles in potassium hydrogen chloromaleate. The atoms may be identified by comparison with Fig. 3. Top diagram: two sets of interatomic distances. The upper number of each pair is the distance (Å) with its least-squares standard error ($\text{\AA} \times 10^{-3}$) in parenthesis, computed directly from the atomic parameters of Table 1. The lower number is the distance (Å) corrected for the effects of thermal motion as described in the text. Numbers with underline are the deviations (Å) from the least-squares best plane through the 11 atoms of the anion, with standard errors ($\text{\AA} \times 10^{-3}$) in parenthesis. Bottom diagram: bond angles (degrees) with least-squares standard errors (degrees $\times 10^{-1}$) in parenthesis.

the *c* axis. The polyhedra are linked into infinite sheets by sharing edges.

Thermal motion

In Figs. 3 and 4 the atoms are represented by the ellipsoids for thermal displacement including 50% probability. The figures suggest that the observed displacements can be accounted for as resulting mainly from rigid body oscillation and libration. An analysis on this basis by the method of Cruickshank (1956) did not yield a highly satisfactory fit, as the motions of some atoms were significantly overestimated and others were underestimated. It is plausible that the disagreement is due to neglect of internal modes of vibration, such as librations about the carbon-carbon single bonds (the largest discrepancies are for displacements out of the plane of the ion). Nevertheless, the analysis yields parameters that are plausible and are probably roughly correct. The principal root-mean-square values of the angular displacements resulting from this calculation (H and K atoms omitted, center of motion at center of mass of anion) are 5.3° , 3.5° and 3.0° . The first principal axis is within 5° of the C(1)-C(4) direction and the second is within 4° of the normal to the plane of the ion. This result is not greatly changed if the hydrogen atoms are included, nor if the center of motion is moved as much as 0.1 \AA from the center of mass.

Corrected bond distances

The libration amplitudes just described were used to estimate thermal corrections to the bond distances, according to the formula (Busing & Levy, 1964)

$$S = S_0 [1 + 0.5 \sum_{i=1}^3 (\overline{\varphi_i^2} \sin^2 \psi_i)]$$

in which S is the estimate of the mean separation, S_0 the uncorrected bond distance, and $\overline{\varphi_i^2}$ a mean-square amplitude of libration about a principal axis i making angle ψ_i with the bond direction. These corrected distances are shown in Fig. 5. In the case of the C(3)-H(2) bond, special treatment is required because the rigid-body model of thermal motion underestimates the thermal amplitudes of H(2), and the discrepancy clearly arises from bending of the C-H bond. Hence, an additional correction to the C-H bond length was made, using these excess thermal amplitudes and the 'riding' model discussed by Busing & Levy (1964). The results are as follows: uncorrected 1.088 \AA , corrected for libration only 1.094 \AA , corrected for both libration and bending 1.111 \AA . We feel the accuracy of the corrected values to be of the order of 0.01 \AA .

The bridging hydrogen atom

We consider now possible interpretations of the apparent central location of atom H(1) along the O...H...O bridge. The neutron diffraction analysis has determined the mean position to be at the center, in the sense of an average position over the lattice as well as over time. We ask whether the component of

mean-square displacement of H(1) in the direction of the O...O bridge, as given by the thermal tensor, is to be ascribed to thermal oscillation about a single equilibrium point, or whether a part of this quantity should be ascribed to the existence of equilibrium positions displaced from the center and occupied at random throughout the lattice. The Bragg reflection data, of course, give no direct information on this question.

We shall discuss this question in the same way as was done by McGaw & Ibers (1963) and Ibers (1964) for the analogous question in the bifluoride ion; that is, we compare the mean-square amplitude of stretching of the O-H links in the bridge as deduced from the diffraction data with that to be expected from spectroscopic considerations. For the former, a satisfactory approximation may be taken to be the difference between the mean-square component of displacement of H(1) and the average of those for O(2) and O(3). For the latter, the mean-square displacement of a harmonic oscillator with the reduced mass of a hydrogen atom and a characteristic frequency $\nu = 1600 \text{ cm}^{-1}$ is chosen. We neglect the effect of oscillations in which O(2) and O(3) move symmetrically and H(1) is stationary, since the effective mass and the rigidity of the ring structure would both tend to make such contributions small. The frequency 1600 cm^{-1} is the upper limit of the region of the infrared spectrum of potassium hydrogen maleate (not chlorinated) which is sensitive to replacement of H by D (Cardwell, Dunitz & Orgel, 1953) and is a value to be attributed to a very strong O...H...O bond (Rundle & Parasol, 1952). We obtain $0.017 \pm 0.002 \text{ \AA}^2$ from the diffraction data *vs.* 0.010 \AA^2 from the expression $\bar{z}^2 = h/(8\pi^2 \nu m)$, given by harmonic oscillator theory. The difference is not a large one and is plausibly accounted for by some departure of the hydrogen atom potential function from a single parabolic minimum. For example, if a double minimum well is assumed, the square of the displacement of an equilibrium point from the center must be of the order of 0.007 \AA^2 , the difference between the above two values (the effects of displaced equilibrium and thermal amplitude are additive in the total mean-square displacement), and is thus less than the predicted thermal mean-square displacement of 0.010 \AA^2 . In other words, an excursion having the r.m.s. magnitude of 0.1 \AA from one of the (assumed) displaced equilibrium points would bring the hydrogen atom to within this distance of the second equilibrium point; this situation would imply a low barrier, probably not greater than the zero-point vibrational level. In view of this comparison, the notion of a static displacement of the equilibrium position, involving a barrier much higher than the zero-point energy, can hardly be maintained. We therefore conclude that the hydrogen atom is effectively centered in a substantially symmetric potential well.

The stretching amplitude of the C-H Bond

Considerations similar to the foregoing may also be applied to the stretching amplitude of the C(3)-H(2)

bond. Under similar assumptions, the diffraction analysis yields $0.0048 \pm 0.0018 \text{ \AA}^2$ as compared with a theoretical value of 0.0055 \AA^2 for $\nu = 3020 \text{ cm}^{-1}$ (Herzberg, 1945). This very satisfactory agreement lends confidence to the reality of the thermal parameters determined by diffraction methods, and to the correctness of their interpretation.

The Fourier syntheses were computed with the program of A. Zalkin of Lawrence Radiation Laboratory. The least-squares refinements utilized a modification of the Busing, Martin & Levy (1962) program, prepared in part by W. C. Hamilton and J. A. Ibers of Brookhaven National Laboratory and by C. K. Johnson of this Laboratory. Distances, angles, and other derived quantities were computed with C. K. Johnson's modification of the Busing, Martin & Levy (1964) function and error program. The best-plane computation utilized a program written by W. C. Hamilton. The rigid-body thermal analysis program was that of Gantzel, Coulter & Trueblood (1961). The program for stereoscopic illustrations was prepared at this Laboratory by C. K. Johnson (1965); that for preparation of Table 2 by G. M. Brown and K. O. Martin of this Laboratory. The crystal specimen was grown by Dr S. W. Peterson, formerly of this Laboratory.

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